- **7.** The order parameter on a Bethe lattice with coordination number z
  - (i)  $P_{\infty}(p)$  is the probability that an arbitrarily selected site belongs to the percolating infinite cluster. Consider the 'origin' in the Bethe lattice.

$$P_{\infty}(p) = \text{probability 'origin' is occupied } \cdot$$
  
probability at least one of the z branches connects to infinity  
$$= p \left[ 1 - Q_{\infty}^{z}(p) \right]$$
(7.1)

where  $Q_{\infty}(p)$  denotes the probability that a given branch does *not* connect to infinity. Again, we will rely on the fact that all sites in a Bethe lattice are equivalent, so  $Q_{\infty}(p)$  is also the probability that a subbranch does not connect to infinity. Hence

$$Q_{\infty}(p) = \text{neighbour to 'origin' is empty} + \text{neighbour to 'origin' is occupied}$$
  
but none of the  $(z - 1)$  subbranches connect to infinity  
$$= (1 - p) + pQ_{\infty}^{z-1}(p)$$
(7.2)

(ii) For convenience, we drop the *p*-dependence of  $Q_{\infty}(p)$  and simply write  $Q_{\infty}$ . Let

$$Q_{\infty}^{z-1} = (1 - [1 - Q_{\infty}])^{z-1} = (1 - x)^{z-1}$$
 with  $x = 1 - Q_{\infty}$ .

We expand to second order in x around x = 0.

$$f(x) = (1-x)^{z-1} \Rightarrow f(0) = 1$$
  

$$f^{(1)}(x) = -(z-1)(1-x)^{z-2} \Rightarrow f^{(1)}(0) = -(z-1)$$
  

$$f^{(2)}(x) = (z-1)(z-2)(1-x)^{z-3} \Rightarrow f^{(2)}(0) = (z-1)(z-2)$$
(7.3)

implying

$$Q_{\infty}^{z-1} \approx f(0) + f^{(1)}(0)x + \frac{1}{2}f^{(0)}(0)x^{2} + \dots$$
  
=  $1 - (z-1)x + \frac{1}{2}(z-1)(z-2)x^{2} + \dots$   
=  $1 - (z-1)(1-Q_{\infty}) + \frac{1}{2}(z-1)(z-2)(1-Q_{\infty})^{2} + \dots$  (7.4)

Using the Taylor expansion result in Equation (7.2) we find

$$Q_{\infty} = 1 - p + pQ_{\infty}^{z-1}$$

$$\approx 1 - p + p - p(z-1)(1 - Q_{\infty}) + p\frac{1}{2}(z-1)(z-2)(1 - Q_{\infty})^{2}$$

$$= 1 - p(z-1) + p(z-1)Q_{\infty} + a(1 + Q_{\infty}^{2} - 2Q_{\infty})$$
(7.5)

and rearranging

$$aQ_{\infty}^{2} + \{ p(z-1) - 1 - 2a \} Q_{\infty} + a + 1 - p(z-1) = 0 \Leftrightarrow aQ_{\infty}^{2} + (b - 2a)Q_{\infty} + a - b = 0 \Leftrightarrow Q_{\infty} = \frac{2a - b \pm \sqrt{(b - 2a)^{2} - 4a(a - b)}}{2a} = \frac{2a - b \pm \sqrt{b^{2}}}{2a}.$$
 (7.6)

As b > 0 since p eventually is larger that  $\frac{1}{z-1}$  we find

$$Q_{\infty} = \begin{cases} 1 & \text{for } p \le p_c \\ \frac{a-b}{a} & \text{for } p > p_c. \end{cases}$$

The solution  $Q_{\infty} = 1 \Rightarrow P_{\infty}(p) = 0$  belongs to the regime  $p \le p_c$ . The other solution is nontrivial and belongs to the regime  $p > p_c$ , and

$$Q_{\infty} = \frac{a-b}{a} = 1 - \frac{2p(z-1)-2}{p(z-1)(z-2)}.$$

(iii) The relevant solution has  $Q_{\infty} < 1$ . Substituting into the Equation (7.1)

$$P_{\infty}(p) = p(1 - Q_{\infty}^{z})$$
  
=  $p\left[1 - \left(1 - \frac{2p(z-1) - 2}{p(z-1)(z-2)}\right)^{z}\right]$   
=  $p\left[1 - \left(1 - \frac{b}{a}\right)^{z}\right]$   
=  $p - p(1 - \frac{b}{a})^{z}.$  (7.7)

Note that the ratio

$$\frac{b}{a} = \frac{p(z-1)-1}{p\frac{1}{2}(z-1)(z-2)} \to 0 \qquad \text{for } p \to \frac{1}{z-1} = p_c,$$

so  $\frac{b}{a}$  is a small quantity for  $p \to p_c$ . Let  $g(x) = (1-x)^z$ . Taylor expanding to first order we find  $g(0) = 1, g^{(1)}(x) = -z(1-x)^{z-1}, g^{(1)}(0) = -z$  so  $(1-x)^z \approx 1 - zx$  for  $x \to 0$ . Thus

$$P_{\infty}(p) = p - p(1 - \frac{b}{a})^{z}$$
  

$$\approx p - p(1 - z\frac{b}{a}) \text{ for } p \to p_{c}$$
  

$$= pz\frac{b}{a}$$
  

$$= pz\frac{2p(z - 1) - 2}{p(z - 1)(z - 2)}$$
  

$$= \frac{2z}{z - 2} \left( p - \frac{1}{z - 1} \right)$$
  

$$= \frac{2z}{z - 2} (p - p_{c})$$
(7.8)

with

$$p_{c} = \frac{1}{z - 1}$$
$$A = \frac{2z}{z - 2} \stackrel{z=3}{=} 6.$$

and



Figure 5.1: (a) The d = 1 cluster number density, n(s, p), as a function of cluster size, s. The five curves correspond to p = 0.4, 0.905, 0.99, 0.999, 0.9999, respectively. (b) The transformed cluster number densities  $s^2n(s, p)$  versus the cluster size s. (c) For each transformed cluster number density  $s^{\tau}n(s, p)$ , the argument is rescaled from s to  $s/s_{\xi}$  where the characteristic cluster size,  $s_{\xi} = -1/\ln(p) \approx 1, 10, 100, 1000, 10000$ , respectively.

- 5. (i) See Figure 5.1 above.
  - (ii) The cluster number density in d = 1 can be written on the form

$$n(s,p) = s^{-2} \mathcal{G}_{1d} \left( s/s_{\xi} \right)$$
(5.1)

where the function

$$\mathcal{G}_{1d}(s/s_{\xi}) = (s/s_{\xi})^2 \exp(-s/s_{\xi}).$$
(5.2)

Thus the transformed cluster number density

$$s^2 n(s,p) = \mathcal{G}_{\mathrm{1d}}\left(s/s_{\xi}\right) \tag{5.3}$$

is only a function of the ratio  $s/s_{\xi}$  and by plotting  $s^2n(s,p)$  versus  $s/s_{\xi}$ , all the graphs collapse onto one curve, the graph of the function  $\mathcal{G}_{1d}$ . Note that

$$\mathcal{G}_{1\mathrm{d}}\left(s/s_{\xi}\right) = \left(s/s_{\xi}\right)^{2} \exp\left(-s/s_{\xi}\right)$$
$$= \begin{cases} \left(s/s_{\xi}\right)^{2} & \text{for } s/s_{\xi} \ll 1\\ \text{decays rapidly} & \text{for } s/s_{\xi} \gg 1, \end{cases}$$
(5.4)

so for small arguments  $s \ll s_{\xi}$ , the function  $\mathcal{G}_{1d}$  increases as the argument squared while it decays exponentially fast for large arguments  $s \gg s_{\xi}$ , see Figure 5.1(c).